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Redetermination of the crystal structure of uracil. By ROBERT F. STEWART, *Mellon Institute, 4400 Fifth Avenue, Pittsburgh, Pennsylvania 15213, U.S.A.* and LYLE H. JENSEN, *Department of Biological Structure, University of Washington, Seattle, Washington 98105, U.S.A.*

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The crystal structure of uracil, originally reported by Parry in 1954, has been redetermined. The unit-cell parameters are $a=11.938 \pm 0.001$, $b=12.376 \pm 0.0009$, $c=3.6552 \pm 0.0003$ Å; $\beta=120^\circ 54' \pm 0.4'$. The space group is $P2_1/a$. Three dimensional X-ray diffraction data were measured by diffractometry with Mo $K\alpha$ radiation. The final R was 0.045. The redetermined structure is essentially a confirmation of Parry's structure. The C(5)-C(6) bond length (1.340 Å, e.s.d. 0.002 Å) in this work, however, is 0.07 Å shorter than the value (1.41 Å, e.s.d. 0.011 Å) reported by Parry.

The crystal structure of uracil was first reported by Parry (1954). One interesting feature of the structure is the close approach of an oxygen atom to carbon atoms in different molecules. The reported distances are 3.19 and 3.28 Å. Another interesting result is the C(5)-C(6) bond length of 1.41 ± 0.01 Å, a value considerably greater (0.07 Å) than for a similar bond found in 1-methylthymine (Hoogsteen, 1963a, b), thymine monohydrate (Gerdil, 1961) and 1-methyluracil (Green, Matthews & Rich, 1962). The accuracy of the results, unfortunately, was limited by the quality of the uracil crystal. We report here the structural results of an X-ray diffraction investigation of a uracil crystal.

High quality single crystals of uracil were grown by heating the powder to its decomposition point in a Pyrex test

tube with a sharp temperature gradient. Single crystals of stubby habit with prominent forms {001} and {110} grew on the walls of the tube at a temperature somewhat below the decomposition point. One precession and several Weissenberg photographs confirmed the space group $P2_1/a$ as reported by Parry (1954). Accurate cell parameters were determined on a Picker X-ray diffractometer with a General Electric quarter circle goniostat. The parameters for the monoclinic cell are:

$$\begin{aligned} a &= 11.938 \pm 0.001 \text{ Å} \\ b &= 12.376 \pm 0.0009 \\ c &= 3.6552 \pm 0.0003 \\ \beta &= 120^\circ 54' \pm 0.4' \\ \lambda(\text{Mo } K\alpha) &= 0.71069 \text{ Å.} \end{aligned}$$

Table 1. Final atomic parameters

The temperature factors are of the form

$$T_i = \exp [-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)] \text{ and } T_H = \exp [-B(\sin \theta/\lambda)^2].$$

Estimated standard deviations, in parenthesis, are right-adjusted:

e.g. For O(7), $\sigma(x)=11 \cdot 10^{-5}$ and for H(9), $\sigma(x)=16 \cdot 10^{-4}$.

	x	y	z	$10^4\beta_{11}$	$10^4\beta_{22}$	$10^4\beta_{33}$	$10^4\beta_{12}$	$10^4\beta_{13}$	$10^4\beta_{23}$
O(7)	0.08497 (11)	0.28677 (9)	-0.02173 (43)	80 (1)	41 (0.8)	1670 (22)	11 (0.8)	252 (4)	30 (3)
O(8)	0.12987 (10)	-0.07706 (8)	0.01962 (44)	66 (1)	35 (0.7)	1537 (19)	-6 (0.8)	228 (4)	-17 (3)
N(1)	0.27283 (12)	0.21750 (11)	0.05336 (50)	58 (1)	36 (0.9)	1157 (22)	-7 (0.9)	176 (4)	5 (4)
N(3)	0.11208 (12)	0.10468 (10)	0.00406 (48)	43 (1)	37 (0.8)	1040 (19)	-2 (0.8)	143 (4)	0 (3)
C(2)	0.15191 (14)	0.20933 (12)	0.00913 (56)	58 (1)	41 (1)	924 (22)	-3 (1)	143 (5)	10 (4)
C(4)	0.17988 (14)	0.01237 (12)	0.03494 (54)	52 (1)	40 (1)	829 (21)	-2 (1)	129 (5)	-10 (4)
C(5)	0.30519 (14)	0.02857 (13)	0.08193 (58)	51 (1)	44 (1)	1026 (25)	2 (1)	158 (5)	-9 (4)
C(6)	0.34620 (15)	0.12977 (12)	0.08905 (59)	49 (1)	51 (1)	940 (22)	-6 (1)	145 (5)	-14 (4)
H(9)	0.3536 (16)	-0.0322 (14)	0.0986 (56)	3.41 (39): B_H					
H(10)	0.4279 (16)	0.1464 (13)	0.1120 (50)	3.04 (35)					
H(11)	0.2964 (15)	0.2804 (14)	0.0416 (52)	3.15 (39)					
H(12)	0.0343 (16)	0.0947 (13)	-0.0320 (52)	3.33 (39)					

The parameters a and c are 1 % greater than those reported by Parry and the angle β is 0.9° larger. The discrepancy evidently lies in the determination of the angle β , since b_{100} and d_{001} for both sets of measurements agree well.

The Bragg scattered intensities were measured for the upper hemisphere (the crystal was mounted along c^* with Mo $K\alpha$ radiation on a Picker diffractometer. The 2θ scan technique was employed (Alexander & Smith, 1964). All

reflections out to $2\theta=50^\circ$ ($\sin \theta/\lambda=0.595 \text{ \AA}^{-1}$) were measured; in the shell from 2θ of 50 to 80°, only reflections with expected net intensities greater than 200 counts were measured. The overall agreement, R_{exp} , for the equivalent reflections hkl and $\bar{h}\bar{k}\bar{l}$ was 0.0168. We define $R_{\text{exp}}=\Sigma |I_{hkl}-I_{\bar{h}\bar{k}\bar{l}}| / \Sigma I_{\text{mean}}$. The most intense reflections were measured by powder diffractometry. These intensities, from the planes (001), (200), and (120), were judged free from

Table 2. List of structure factors for uracil

The three columns of each row contain the values for k , $10F_0$, $10F_c$. The F_c values were calculated from the parameters in Table 1.

	$0, k, 0$	$6, 10, 0$	$1, 3, k+1$	$7, 202, -205$	$-16, k, 1$	$14, 21, -230$	$12, 66, -57$	$10, 25, -26$	$12, 42, -38$	$3, 35, -37$
2	55	-51	210	36	1	98	55	62	1	42
2	201	-124	125	42	-46	2	27	-77	2	43
2	120	-118	125	43	-37	3	32	-73	3	43
10	210	-212	124	45	-37	4	36	-73	4	50
10	100	-101	100	45	-37	5	35	-73	5	48
16	69	+66	88	84	-82	6	30	-70	6	48
1	1, k, 0	0	10	12	-10	7	59	-55	7	37
1	216	-231	70	70	-77	8	58	-52	8	37
1	140	-125	42	42	-40	9	50	-46	9	37
1	120	-125	67	18	-28	10	42	-40	10	37
5	57	-52	84	17	-24	11	39	-38	11	37
5	54	-52	80	34	-36	12	39	-38	12	37
10	110	-111	115	36	-34	13	39	-38	13	37
10	70	-72	111	35	-34	14	39	-38	14	37
16	17	-13	15	15	-8	15	39	-38	15	37
2	2, k, 0	0	10	10	-10	16	50	-46	16	37
2	216	-231	70	70	-77	17	50	-46	17	37
2	140	-125	42	42	-40	18	50	-46	18	37
2	120	-125	67	18	-28	19	50	-46	19	37
5	57	-52	84	17	-24	20	50	-46	20	37
5	54	-52	80	34	-36	21	50	-46	21	37
10	110	-111	115	36	-34	22	50	-46	22	37
10	70	-72	111	35	-34	23	50	-46	23	37
16	17	-13	15	15	-8	24	50	-46	24	37
2	3, k, 0	0	10	10	-10	25	50	-46	25	37
2	216	-231	70	70	-77	26	50	-46	26	37
2	140	-125	42	42	-40	27	50	-46	27	37
2	120	-125	67	18	-28	28	50	-46	28	37
5	57	-52	84	17	-24	29	50	-46	29	37
5	54	-52	80	34	-36	30	50	-46	30	37
10	110	-111	115	36	-34	31	50	-46	31	37
10	70	-72	111	35	-34	32	50	-46	32	37
16	17	-13	15	15	-8	33	50	-46	33	37
2	4, k, 0	0	11	0	-11	33	50	-46	33	37
2	216	-231	70	70	-77	34	50	-46	34	37
2	140	-125	42	42	-40	35	50	-46	35	37
2	120	-125	67	18	-28	36	50	-46	36	37
5	57	-52	84	17	-24	37	50	-46	37	37
5	54	-52	80	34	-36	38	50	-46	38	37
10	110	-111	115	36	-34	39	50	-46	39	37
10	70	-72	111	35	-34	40	50	-46	40	37
16	17	-13	15	15	-8	41	50	-46	41	37
2	5, k, 0	0	12	0	-12	41	50	-46	41	37
2	216	-231	70	70	-77	42	50	-46	42	37
2	140	-125	42	42	-40	43	50	-46	43	37
2	120	-125	67	18	-28	44	50	-46	44	37
5	57	-52	84	17	-24	45	50	-46	45	37
5	54	-52	80	34	-36	46	50	-46	46	37
10	110	-111	115	36	-34	47	50	-46	47	37
10	70	-72	111	35	-34	48	50	-46	48	37
16	17	-13	15	15	-8	49	50	-46	49	37
2	6, k, 0	1	36	32	-34	50	50	-46	50	37
2	216	-231	70	70	-77	51	50	-46	51	37
2	140	-125	42	42	-40	52	50	-46	52	37
2	120	-125	67	18	-28	53	50	-46	53	37
5	57	-52	84	17	-24	54	50	-46	54	37
5	54	-52	80	34	-36	55	50	-46	55	37
10	110	-111	115	36	-34	56	50	-46	56	37
10	70	-72	111	35	-34	57	50	-46	57	37
16	17	-13	15	15	-8	58	50	-46	58	37
2	7, k, 0	0	13	32	-34	59	50	-46	59	37
2	216	-231	70	70	-77	60	50	-46	60	37
2	140	-125	42	42	-40	61	50	-46	61	37
2	120	-125	67	18	-28	62	50	-46	62	37
5	57	-52	84	17	-24	63	50	-46	63	37
5	54	-52	80	34	-36	64	50	-46	64	37
10	110	-111	115	36	-34	65	50	-46	65	37
10	70	-72	111	35	-34	66	50	-46	66	37
16	17	-13	15	15	-8	67	50	-46	67	37
2	8, k, 0	0	14	32	-34	68	50	-46	68	37
2	216	-231	70	70	-77	69	50	-46	69	37
2	140	-125	42	42	-40	70	50	-46	70	37
2	120	-125	67	18	-28	71	50	-46	71	37
5	57	-52	84	17	-24	72	50	-46	72	37
5	54	-52	80	34	-36	73	50	-46	73	37
10	110	-111	115	36	-34	74	50	-46	74	37
10	70	-72	111	35	-34	75	50	-46	75	37
16	17	-13	15	15	-8	76	50	-46	76	37
2	9, k, 0	0	15	32	-34	77	50	-46	77	37
2	216	-231	70	70	-77	78	50	-46	78	37
2	140	-125	42	42	-40	79	50	-46	79	37
2	120	-125	67	18	-28	80	50	-46	80	37
5	57	-52	84	17	-24	81	50	-46	81	37
5	54	-52	80	34	-36	82	50	-46	82	37
10	110	-111	115	36	-34	83	50	-46	83	37
10	70	-72	111	35	-34	84	50	-46	84	37
16	17	-13	15	15	-8	85	50	-46	85	37
2	10, k, 0	0	16	32	-34	86	50	-46	86	37
2	216	-231	70	70	-77	87	50	-46	87	37
2	140	-125	42	42	-40	88	50	-46	88	37
2	120	-125	67	18	-28	89	50	-46	89	37
5	57	-52	84	17	-24	90	50	-46	90	37
5	54	-52	80	34	-36	91	50	-46	91	37
10	110	-111	115	36	-34	92	50	-46	92	37
10	70	-72	111	35	-34	93	50	-46	93	37
16	17	-13	15	15	-8	94	50	-46	94	37
2	11, k, 0	0	17	32	-34	95	50	-46	95	37
2	216	-231	70	70	-77	96	50	-46	96	37
2	140	-125	42	42	-40	97	50	-46	97	37
2	120	-125	67	18	-28	98	50	-46	98	37
5	57	-52	84	17	-24	99	50	-46	99	37
5	54	-52	80	34	-36	100	50	-46	100	37
10	110	-111	115	36	-34	101	50	-46	101	37
10	70	-72	111	35	-34	102	50	-46	102	37
16	17	-13	15	15	-8	103	50	-46	103	37
2	12, k, 0	0	18	32	-34	104	50	-46	104	37
2	216	-231	70	70	-77	105	50	-46	105	37
2	140	-125	42	42	-40	106	50	-46	106	37
2	120	-125	67	18	-28	107	50	-46	107	37
5	57	-52	84	17	-24	108	50	-46	108	37
5	54	-52	80	34	-36	109	50	-46	109	37
10	110	-111	115	36	-34	110	50	-46	110	37
10	70	-72	111	35	-34	111	50	-46	111	37
16	17	-13	15	15	-8	112	50	-46	112	37
2	13, k, 0	0	19	32	-34	113	50	-46	113	37
2	216	-231	70	70	-77	114	50	-46	114	37
2	140	-125	42	42	-40	115	50	-46	115	37
2	120	-125	67	18	-28	116	50	-46	116	37
5	57	-52	84	17	-24	117	50	-46	117	37
5	54	-52	80	34	-36	118	50	-46	118	37
10	110	-111	115	36	-34	119	50	-46	119	37
10	70	-72	111	35	-34	120	50	-46	120	37
16	17	-13	15	15	-8	121	50	-46	121	37
2	14, k, 0	0	20	32	-34	122	50	-46	122	37
2	216	-231	70	70	-77	123	50	-46	123	37
2	140	-125	42	42	-40	124	50	-46	124	37
2	120	-125	67	18	-28	125	50	-46	125	37
5	57	-52	84	17	-24	126	50	-46</td		

secondary extinction. Adsorption differences were found to be small ($\sim 1.5\%$) and no corrections were made.

The averaged intensities were converted to structure factors within the ideal mosaic approximation and with appropriate account of Lorentz and polarization effects. Weights were based on the counting statistics and on other fluctuations in the diffractometer – the latter amounting to a 1 % relative standard error. A total of 1163 unique reflections, of which 132 were unobserved, were used in the least-squares analysis.

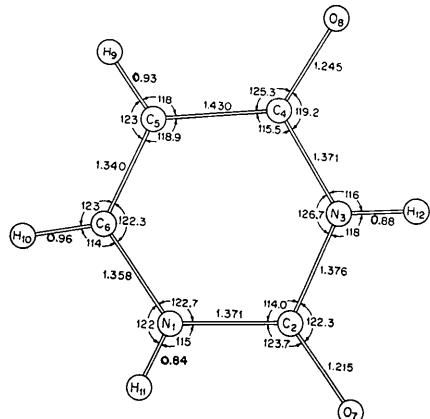


Fig. 1. Dimensions of the uracil molecule.

The structure factor calculation with Parry's (1954) atomic and thermal parameters for all atoms yielded an R of 0.18. The scattering factors for O, N, and C were from Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955) and for H were from Stewart, Davidson & Simpson (1965). Refinement proceeded by full-matrix least-squares whereby an overall scale factor, atomic positions, and anisotropic thermal parameters (isotropic for H atoms) were varied. For these 89 parameters the sum of the weighted, squared residuals, $w(|F_0| - |F_c|)^2$ was minimized. After three cycles the refinement terminated, with marginal shifts in the parameters, the largest being 0.25σ . The final R was 0.045, omitting the unobserved reflections. The final weighted R was 0.037. The goodness of fit was 2.34. The average estimated standard deviation for positions of the heavy atoms was 0.0016 Å; for the hydrogen 0.018 Å. The final atomic parameters are listed in Table 1. All calculations were carried out on an IBM 7094 computer with programs from X-RAY 63 (Stewart *et al.*, 1964).

The observed and calculated structure factors are listed in Table 2.

The distances and angles among the least-squares adjusted precalculated scattering centers are listed in Table 3 and shown in Fig. 1. For comparison the values reported by Parry (1954), corrected to our unit cell parameters are also shown in Table 3. The C(5)-C(6) bond is 0.07 Å shorter than Parry's value, but is in agreement with the C(5)-C(6) bond in thymine (Gerdil, 1961) and 1-methylthymine (Hoog-

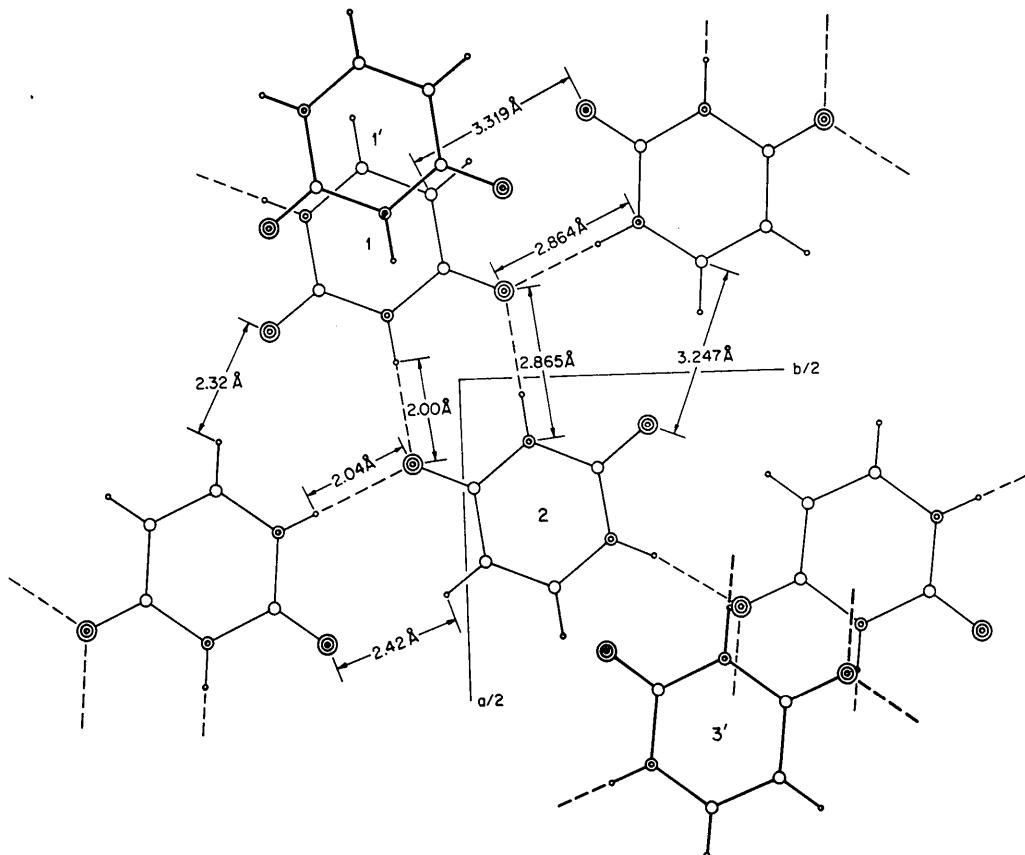


Fig. 2. Orthographic projection of uracil molecules onto (001). Heavy print, molecules near $z=1$. Light print, molecules near $z=0$. Single large circles, carbon atoms. Single small circles, hydrogen atoms. Double circles, nitrogen atoms. Triple circles, oxygen atoms.

steen, 1963a, b). The corresponding bond lengths and angles for 1-methylthymine, also listed in Table 3, are in remarkable agreement with the refined uracil structure. The C(2)-O(7) bond length has not been corrected for the bending motion implied by the thermal parameters. A significantly different bond length is the C(6)-N(1) distance where the bond for uracil (1.358 Å) is 0.025 Å shorter than for 1-methylthymine (1.383 Å). In thymine (Gerdil, 1961), moreover, the C(6)-N(1) bond length is also the same as for 1-methylthymine.

Table 3. Bond distances and angles of uracil

U (1966): this investigation.
U (1954): uracil (Parry, 1954); (corrected to unit-cell parameters in this work).

1-MT: 1-methylthymine (Hoogsteen, 1963a).

Estimated standard deviations in parenthesis

	U (1966)	U (1954)	1-MT
N(1)-C(2)	1.371 (0.002 ₅)	1.354 (0.012)	1.379 (0.004)
C(2)-N(3)	1.376 (0.002 ₀)	1.389 (0.008)	1.379 (0.004)
N(3)-C(4)	1.371 (0.002 ₁)	1.379 (0.010)	1.375 (0.004)
C(4)-C(5)	1.430 (0.002 ₆)	1.421 (0.014)	1.432 (0.004)
C(5)-C(6)	1.340 (0.002 ₃)	1.413 (0.011)	1.346 (0.004)
C(6)-N(1)	1.358 (0.002 ₂)	1.347 (0.012)	1.383 (0.004)
C(2)-O(7)	1.215 (0.002 ₁)	1.236 (0.010)	1.214 (0.004)
C(4)-O(8)	1.245 (0.001 ₉)	1.245 (0.010)	1.237 (0.004)
C(5)-H(9)	0.931 (0.020)		
C(6)-H(10)	0.957 (0.020)		
N(1)-H(11)	0.836 (0.018)		
N(3)-H(12)	0.877 (0.021)		
N(1)-C(2)-N(3)	114.0 (0.1)	115.7	115.4° (0.2)
C(2)-N(3)-C(4)	126.7 (0.2)	124.2	126.3 (0.2)
N(3)-C(4)-C(5)	115.5 (0.1)	118.3	116.1 (0.2)
C(4)-C(5)-C(6)	118.9 (0.2)	116.6	118.3 (0.2)
C(5)-C(6)-N(1)	122.3 (0.2)	121.4	123.3 (0.2)
C(6)-N(1)-C(2)	122.7 (0.1)	123.7	120.6 (0.2)
N(1)-C(2)-O(7)	123.7 (0.2)	124.7	123.3 (0.2)
N(3)-C(2)-O(7)	122.3 (0.2)	119.5	121.3 (0.2)
N(3)-C(4)-O(8)	119.2 (0.2)	118.0	120.0 (0.2)
C(5)-C(4)-O(8)	125.3 (0.2)	123.7	123.9 (0.2)
C(4)-C(5)-H(9)	118.1 (1.2)		
C(6)-C(5)-H(9)	123.0 (1.2)		
C(5)-C(6)-H(10)	123.2 (1.0)		
N(1)-C(6)-H(10)	114.5 (1.0)		
C(6)-N(1)-H(11)	122.12 (1.4)		
C(2)-N(1)-H(11)	115.1 (1.3)		
C(2)-N(3)-H(12)	117.8 (1.0)		
C(4)-N(3)-H(12)	115.5 (1.0)		

The least-squares plane through the eight non-hydrogen atoms of the uracil molecule is

$$-1.2601x + 0.2390y + 3.3162z = -0.1100$$

(x , y , and z in fractions of unit-cell axes) and leads to an average standard deviation of 0.0046 Å. The root mean square average for the estimated standard deviations of these atoms along the direction normal to the plane is 0.0017 Å. On the average the deviation of the atoms from the plane is 2.7σ ; one may conclude that they depart significantly from a plane. The largest departures are for N(3) (-0.0073 Å) and for O(8) (+0.0070 Å). The other distances in Å are 0.0006, 0.0048, 0.0011, -0.0022, -0.0040, and

-0.0001 for atoms O(7), N(1), C(2), C(4), C(5), and C(6), respectively. The least-squares plane for the pyrimidine ring alone [O(7) and O(8) omitted] is given by the equation

$$-1.2404x + 0.2728y + 3.3135z = -0.1004.$$

The standard deviation of the six atoms from this plane is 0.0025 Å. Atoms N(3) and C(4) are appreciably out of this plane; -0.003₄ and +0.003₅ Å, respectively. The atom O(8) is 0.016₇ Å above the plane. The uracil least-squares plane forms a dihedral angle with (001) of 6°9'; the pyrimidine ring least-squares plane is at a similar angle of 6°6'.

Fig. 2 is an orthographic projection of the atoms in the uracil molecules onto (001). The heavy line molecules are at $z=1$; the light line at $z=0$. The N(3)-O(8) hydrogen bond length to the centric related molecule is 2.865 ± 0.002 Å and the N(1)-O(8) length to the screw axis related molecule is 2.864 ± 0.002 Å. These distances are in agreement with Parry's structure (corrected to the a axis reported here.) The corresponding distances from O(8) to H(12) and to H(11) are 2.00 ± 0.02 and 2.04 ± 0.02 Å, respectively. The distance from C(5) to O(7) of a screw related molecule is 3.319 Å; the corresponding H(9)-O(7) length is 2.42 Å. The distance from C(6) to O(7) of a glide related molecule is 3.246 Å and the corresponding H(10)-O(7) distance is 2.32 Å. These approach distances are also in agreement with Parry (again corrected to the present a -axis value). Although the H(10)-O(7) approach, 2.32 Å, is close, it is not as close as the distance reported for the 1-methylthymine crystal structure (Hoogsteen, 1963). For the latter case the O(7)-C(6) distance was 3.11 Å and for O(7)-H(9) the length was 2.14 Å.

Aside from the N(1)-O(8) and N(3)-O(8) hydrogen bonds, the closest heavy atom interatomic distance is 3.20 Å between C(6) of molecule 2 and O'(7) of molecule 3' as depicted in Fig. 2. The other close approaches are C(2')-N(1) 3.317 Å, and C(4')-C(5) 3.323 Å between molecules 1' and 1 in Fig. 2. The spacing between the (001) planes is 3.136 Å.

The details of data gathering and refinement in this investigation will be submitted elsewhere.

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